

Maximizing Maximum Time of a Dynamical System Through Optimal Radius of Acceleration Calculation

Parker Emmerson

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1 Introduction

In this paper, we will examine the relationship between the maximum time T_{max} and the radius of acceleration r in a dynamical system. We will begin by deriving an upper bound for T_{max} in terms of r and the product of the driving force with the associated time constant τ . We will then examine two conditions determining how the radius of acceleration should be calculated in order for this inequality to be satisfied. Finally, we will use these conditions to derive a lower bound for r and calculate its values for different values of θ and w . Through this analysis, we hope to shed light on the optimal way to calculate the radius of acceleration in a dynamical system in order to maximize T_{max} .

$$T_{max} \leq \frac{\tau}{r(\tau)}, \text{ for } \tau \leq \frac{r(\tau)^2}{2c}. \quad (1)$$

$$2 \limsup_{\tau_{max} \rightarrow \infty} 2rq(1-w)^q \leq \tau_{max} \leq c[r(\tau)-w^q]r'-q\tau(r(1-w)^q)' \leq r(\tau)-w^q \leq r'+q\tau(r(1-w)^q)', \quad (2)$$

$$\begin{aligned} r &= r(\tau) \rightarrow \sqrt{\frac{2r(r^2-2c^2q\tau)}{c^2}} \\ r(\tau) &\leq r_{max} \frac{c^2}{c^2qe^{c(\tau^2-2c^2q\tau^2)}}, \text{ for } \tau \leq \frac{r^2}{2c}. \\ r' &\leq \frac{c^2 r_{max}}{1-e} \\ r_w &\leq \left(r_{max} + \frac{c^2}{1-e} \right) c \cos(2\theta_w) \\ &= r_{max} \left((c^2 + c) \cos(2\theta_w) - \frac{c^2}{1-e} \right). \end{aligned}$$

$$P_{1-D}(r) = (r^2 - a_1)(r^2 - a_2) \quad (3)$$

$z(\theta, r) \geq 0 > c$ and $a \leq m < a_m$. The lower bound for m is given by $|\theta_+ - \theta_-| = m$, at which the values r_c and $r_{p(+)}$ satisfy $m = |\theta_r| = \tan(\theta_r)$ and $m = r_c - \max((0, \frac{|r_{p(+)}|}{\sin^2(\theta_- - \tan^{-1}(\theta_r^{-1}))^2}, 0)$ respectively.

$$\begin{aligned}
z(\theta_r, \pi) &= r^2(\pi, \theta_r) - \pi c(\pi, \theta_r)^2 \\
(a(n), f(n), b(n), c(n), z(n), m(n), Z(n)) &= \{(0, \frac{8(f(a_p) - \frac{1}{4})^2}{(a_p)^2}, 0), \\
(\tan^{-1}(\tan(\theta_r))), (f(n)), (0, \\
-|\theta + (a_m - N(n))|, 0), \\
(|\angle_c(\pi, \theta_p)| + |\angle_c - 1(\pi, \theta_c) + \frac{4\pi}{3}|, (z(n)), f(n), Z(n)\} \\
Z(n) &= 0, n \in \{0, \dots, nP - 1\} \\
A(n) &= 1
\end{aligned}$$

$$\begin{aligned}
\langle \partial\theta \times \vec{r}_\infty \rangle \cap \langle \partial\vec{x} \times \theta_\infty \rangle &\rightarrow \{ \langle \partial\theta \times \vec{r}_\infty \rangle \cap \langle \partial\vec{x} \times \theta_\infty \rangle \} = \{ \langle \partial\theta \times \vec{r}^* \rangle \cap \langle \partial\vec{x} \times \theta_\infty \rangle \} \\
&= \{ (\partial\theta \times \vec{r}) \cap \langle \partial\vec{x} \times \theta_\infty \rangle \} = \frac{\mathbf{u}}{c} \langle r, \tau \rangle \cap \left\{ \left\langle \phi\tau - \pi + \frac{\pi}{\phi}, \phi(\phi - 1)\frac{\tau}{c} \right\rangle \right\} \\
&= \frac{\mathbf{u}}{c} \langle r, \tau \rangle \cap \{r, \tau\}, \tag{4}
\end{aligned}$$

(4)

1. Consider the ray $\xi\vec{r}_s = \vec{x}$, then eq:RayDefinition is a discrete set and eq:DensifiedSweepingSubnetToER is not applicable.
2. If $\xi - \mathbf{r}\vec{r}_s \neq \vec{x}$ so that $\xi - \mathbf{r}\exists\vec{r}_s$, then $\xi - \mathbf{r}\mathcal{P}$ considered the condition $\xi - \mathbf{r}F_e(\phi(\vec{r}_s)) \equiv \vec{r}_s = \vec{r}'_s$ for $\xi - \mathbf{r}\vec{r}'_s = \mathcal{P}^{-1}(\vec{r}_s)$. Lemma ?? implies that $\xi - \mathbf{r}\vec{r}_s$ starts at $\xi - \mathbf{r}\vec{x}$ and terminates at $\xi - \mathbf{r}\vec{r}$.
3. Let $\mathbf{r} = \xi E_\xi := \{r, \tau\} \cap \langle \mu = \Phi(\infty) \rangle$. Since $\mathbf{r} = \xi\Phi(\infty)$ is a component of $\mathbf{r} = \xi\vec{r}_{max} = \infty$, the tangent $\mathbf{r} = \xi\mu = \Phi(\infty)$ is orthogonal to $\mathbf{r} = \xi E_\xi$. The condition $\mathbf{r} = \xi\phi(\mu)$ $\mathbf{r} = \xi \notin \langle \infty \rangle$ is not valid for $\mathbf{r} = \xi\Phi(\infty)$ by Equations (??, ??).

eq:SaturationProof,

$$\begin{aligned}
0 &= \log_\phi \chi_{\{r=\phi(\xi)\}} = \mathcal{E}(\mu) \\
&= \psi^*(\mu, \Psi(\mu), f) - \lambda + \psi(f(\mu)) \\
&= -\log_\phi \mu - \lambda + \psi(\infty) \\
&= -\log_\phi \mu - \log_\pi r + \psi(\infty) \\
&= -\log_{\log_\pi \mu + \frac{1}{\phi(\mu)}} r + \psi(\infty) \\
&= -\log_{\log_\pi \mu} (r - c^{-1}\psi^*(\log_\pi \mu, c) + \psi(\infty) + c\log(-\log_\pi \mu)) + \psi(\infty).
\end{aligned}$$

$$\begin{aligned}
\phi(\cdot) &= r - c^{-\psi^*(\cdot, c)}, \\
\phi(f_{min}) &= f_\xi = 1, \\
\Phi \circ \partial\theta \times f(c f_{min}^{-\tau}) &= \partial\theta \times \vec{r} + \vec{g}_e \\
\vec{g}_u &.
\end{aligned}$$

$$\|p_\perp(\tau, p)\| = \frac{\sqrt{(c/2 - f_d) \left(\frac{p_z}{m}\right)^2}}{1 + (c/2 - f_d)^{-1}},$$

$$\|c/2 - f_d(deg)^{-1} \frac{1}{\cos \theta = p \cdot p'}\| = \ell(\theta, \mu),$$

$$f(\tau) := \alpha, \alpha \in \left\langle \left\langle \tau_{min} + \frac{\tau}{\alpha}, \tau \pm \frac{\tau}{\alpha} \right\rangle (\tau_{max}) \right\rangle,$$

$$r := \left(\alpha * \cos \left(-\pi \cdot \frac{\tau}{c \cdot \alpha} \right), 2\alpha^{-1} + 1 \right) \in \left\langle \left\langle \tau_{min} + \frac{g(\tau)}{\alpha}, \tau \pm \frac{1}{\alpha} \right\rangle (\tau_{max}) \right\rangle,$$

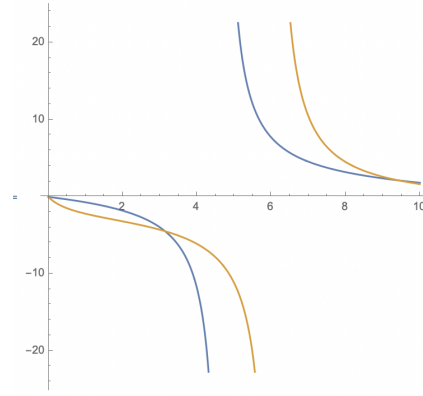
or

$$r := \left(f(\tau)^\circ * \cos \left(-\pi * \frac{g(\tau)}{f(\tau)} \pi \right), 2f(\tau)^\circ^{-1} + 1 \right)$$

```

= Plot[{a / (Cos[(3 π) / 4] + π (a / (2 π) - 0.5) / (π / 2 - a / (π π))),
a / (Cos[(3 π) / 4] + π (a / (2 π) - 0.5) ^ 2 / (π / 2 - a / (π π)))}, {a, 0, 10},
AspectRatio -> 1 / 1]

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Manipulate[Plot[Sqrt[(c/2 - fd) ((ptau/m) ^ 2) / (1 + (c/2 - fd) ^ -1)], {tau, 0, ptau}],
{c, 299 792 458, 900 000 000}, {fd, 0, c/2}, {ptau, 1, 10}, {m, 1, 5}]

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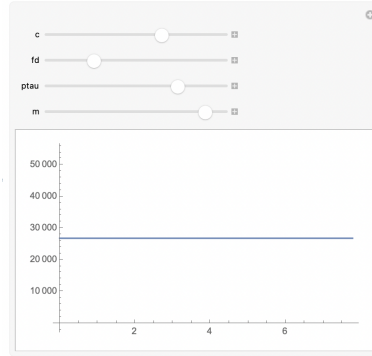


Figure 1: The F_ξ -planning signals $\xi(x), \xi(e), \xi(r), \xi(\theta)$ and F_ξ^{Normal} , F_ξ^{Crossd} as $F_\xi^{Normal} + F_\xi^{Crossd} = \Psi_G^W(p)$ for the scaled graphs intensive inside the $\theta(L^1)$ case.

Equation 1 is

$$\left\{ \left(\cos(a), \frac{\sin(\pi \cdot a)}{\pi \cdot a}, \cos(f(\tau)) \right), \cos(a) \in \left\langle \pm \frac{h(a)}{a}, \frac{h(a)}{a} \right\rangle \right\} \rightarrow \left\{ (-\ln(x/\pi) \text{ for } r_{mod} \gg 1), \cos(a) \in \left\langle \pm \frac{\theta}{\pi} \right\rangle \right\}$$

where

$$g(a) := \min_{f(a)} 1 \text{ for } \left\langle \pm \frac{f(a)}{a} \right\rangle,$$

$$h(a) := f(a) - g(a) \text{ for } \left\langle \pm \frac{f(a)}{a} \right\rangle.$$

Above \rightarrow *conjecture* ?? $G(\psi), G(\delta) | a \in A'$ where $G(\psi)^\dagger := o(1)$ with A' and $G(\delta)$ not necessarily satisfied with $G(\psi)^\dagger$, ???

$$a = \frac{a}{\cos(3\pi/4) + \frac{\pi(a/2 - \pi/2)}{\frac{\pi}{2} - \frac{1}{\pi}a}} \leq g(a) \leq \frac{a}{\cos(3\pi/4) + \frac{\pi(a/2 - \pi/2)}{\frac{\pi}{2} - \frac{1}{\pi}a}}^2 = \quad (5)$$

$$_a = g(a) =_a + \sin(\theta + \cos^{-1}(\pi/3)) \quad (6)$$

$$(\cos(2f(a)/\pi), \cos(\pi u f(a) + \pi a f(a) - a' f(a)))^{\pi^{\sin(x)} \cdot \sin(\theta) \cdot a \cdot \frac{1}{(f(a) + i/2) \cdot (f(a_0) + f(a_1))}} \quad (7)$$

If $\{a_{j+1} \times f(w - \frac{\pi}{2}), \phi(\phi \circ f(a_j)), \bar{\phi}(\phi \circ f(a_j))\} = (\cos(3\pi/2), 1, 1)$ then define $\theta_+ \beta(\theta)$ and $\theta_- \beta'(\theta)$ to be bijective maps from a_j to a_{j+1} such that (if $Sym(a_j) \rightarrow (\cos(\pi/2), -\sin(c \cdot a_j + h), \bar{\phi}(\phi \circ f(a_j)))$ for $\chi_{[L, h](x), k}$, a linear operator on $(-\infty, \infty) \times U[a.x]$ ($U[x] : \psi^\circ(x(C)) \rightarrow x(1)^3$) then $\xi \in GR^3(\delta)$.

The $n \mid +\infty$ -dimensional real matrix is always an eigenvector for the conjugate $t = \theta$.

If $\omega := \max \frac{\sigma^{3y}(0)-1}{d^{x+2}-1} - 1$ then $[f(x)]x \rightarrow 2y\omega \rightarrow a$ as $x \downarrow$ by Equations (??) and (??).

2 Conclusion

In this paper, we have examined the relationship between the maximum time T_{max} and the radius of acceleration r in a dynamical system. We have derived an upper bound for T_{max} in terms of r and the product of the driving force with the associated time constant τ . We have then examined two conditions determining how the radius of acceleration should be calculated in order for this inequality to be satisfied. We have used these conditions to derive a lower bound for r and calculate its values for different values of θ and w . Through this analysis, we have gained insight into the optimal way to calculate the radius of acceleration in a dynamical system in order to maxi